#### **Astronomical Optics**

#### 4.4. Inteferometry on a single aperture

#### **OUTLINE:**

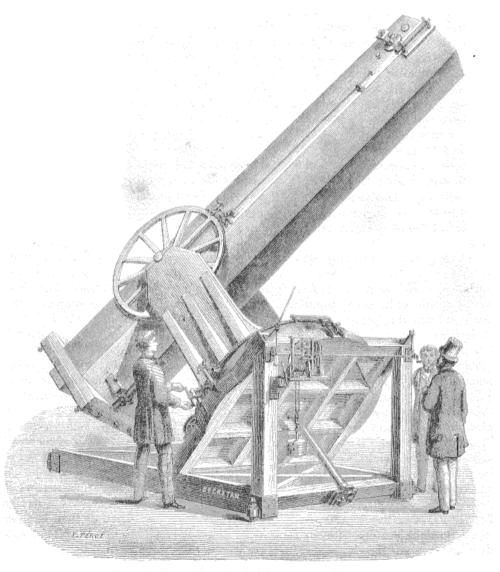
Why interferometry on a single aperture? advantage of interferometric techniques on single aperture telescopes: high precision measurements enabled by good calibration

Aperture masking (slides adapted from presentation by Frantz Martinache)

Pupil remapping

Interferometry as a technique to analyze single aperture short exposures: Speckle interferometry

# Interferometry on a single aperture: First Aperture Masking experiment



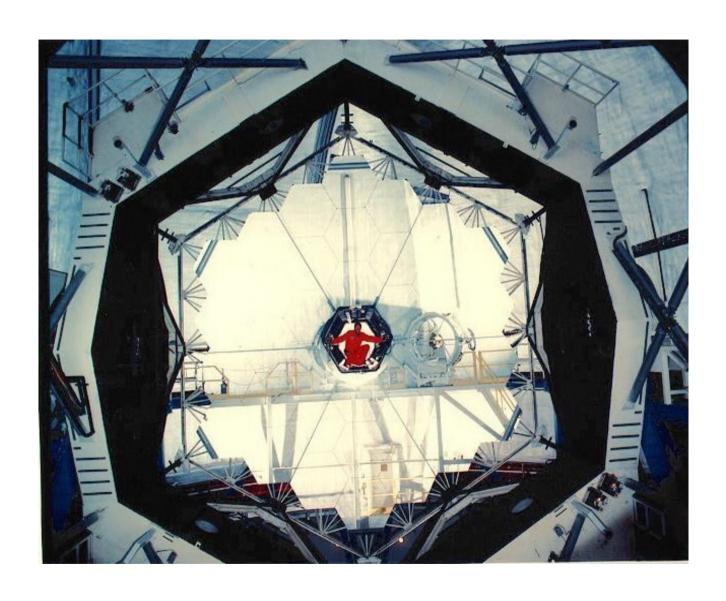
Le grand télescope Foucault, de l'Observatoire de Marseille.

Marseille 1873, Edouard Stephan attempts at measuring the diameter of stars.

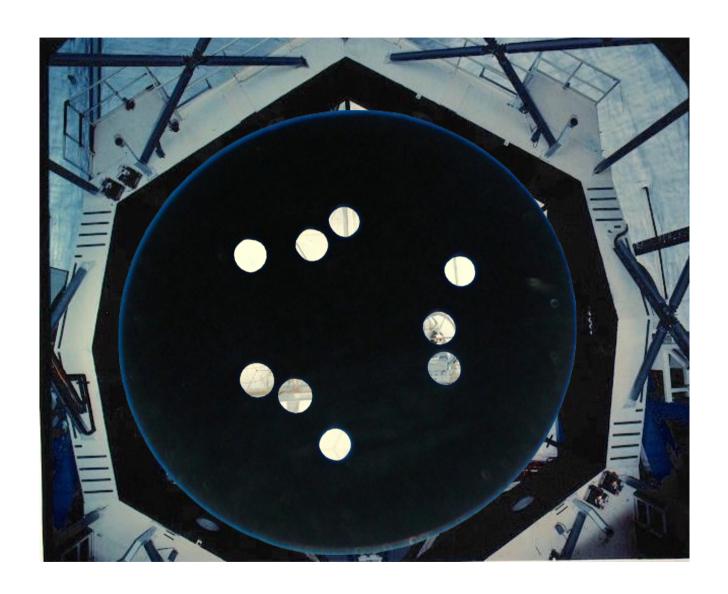
Baseline of 80 cm,  $\emptyset_{\star} < 0.16$ "

Michelson later improved the experiment with an beam expanding the baseline, and was the first to resolve stars

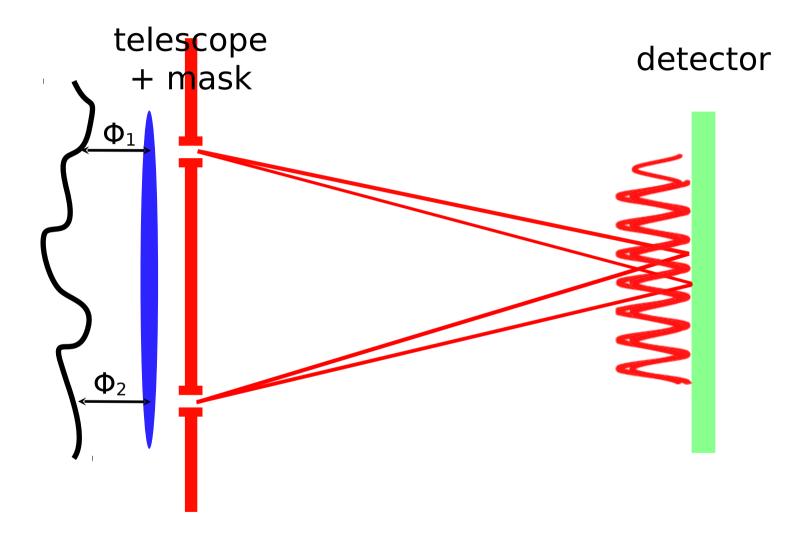
# Aperture masking: principle



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### **Fizeau Interferometry**



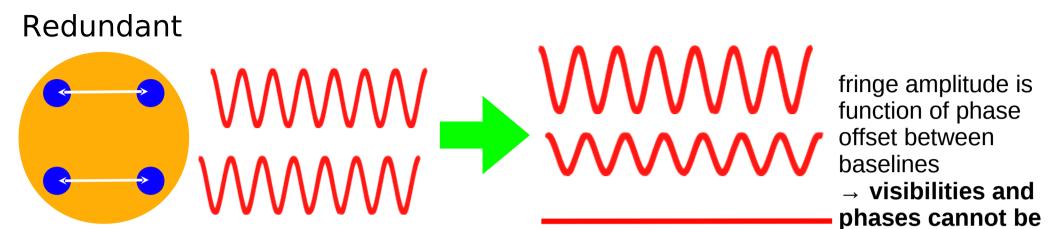
visibility:

$$0 < V^2 < 1$$

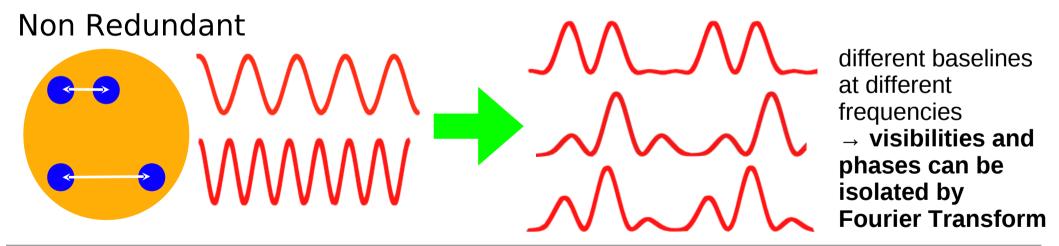
phase:

$$\Phi = \Phi_0 + (\Phi_1 - \Phi_2)$$

# Redundancy: Atmosphere affects the phases, Redundancy destroys the amplitudes



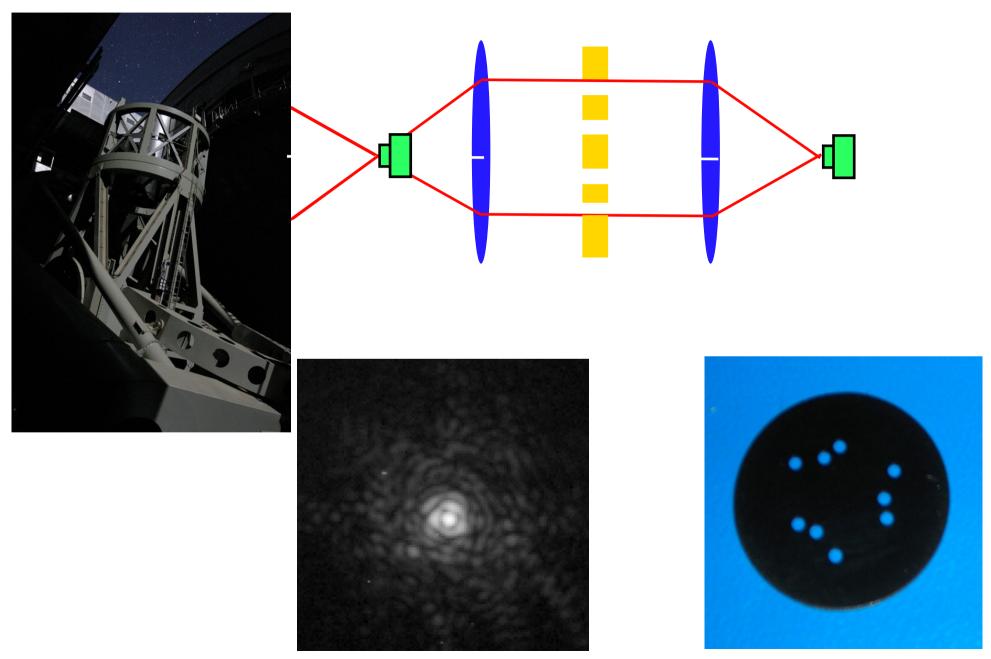
measured





A full aperture is very redundant

## Aperture Masking: creating nonredundant aperture with pupil mask

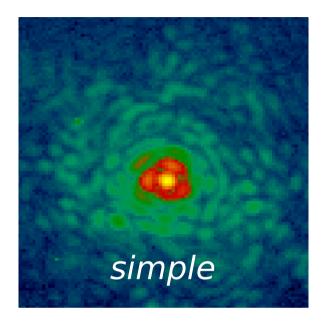


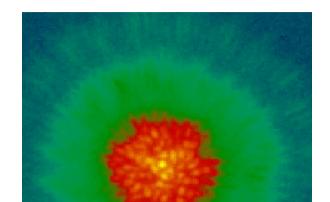
# Image and Fourier planes

conventional imaging

aperture masking

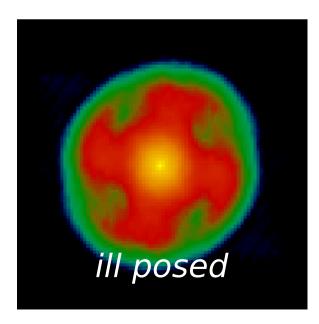


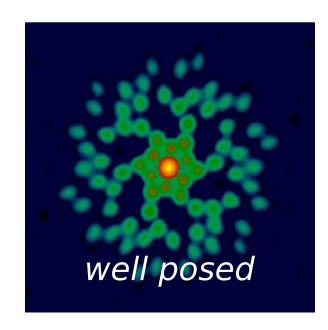




complex

**FT** | 2

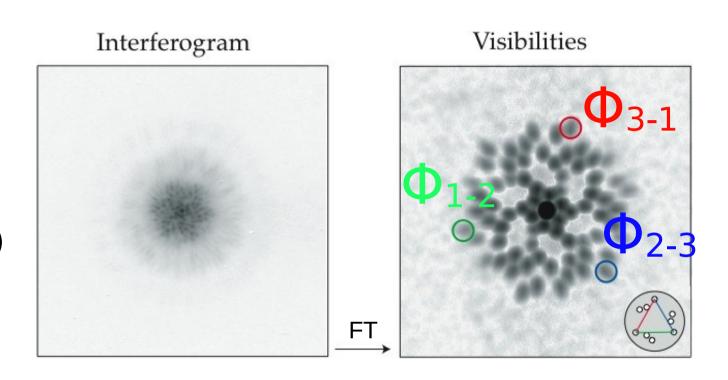




Modulation Transfer Function

### A neat trick: the closure phase

9-hole mask 36 visibilities 84 triangles (28 independent)



$$\Phi(1-2) = \Phi(1-2)_0 + (\Phi_1-\Phi_2)$$

$$\Phi(2-3) = \Phi(2-3)_0 + (\Phi_2-\Phi_3)$$

$$\Phi(3-1) = \Phi(3-1)_0 + (\Phi_3-\Phi_1)$$

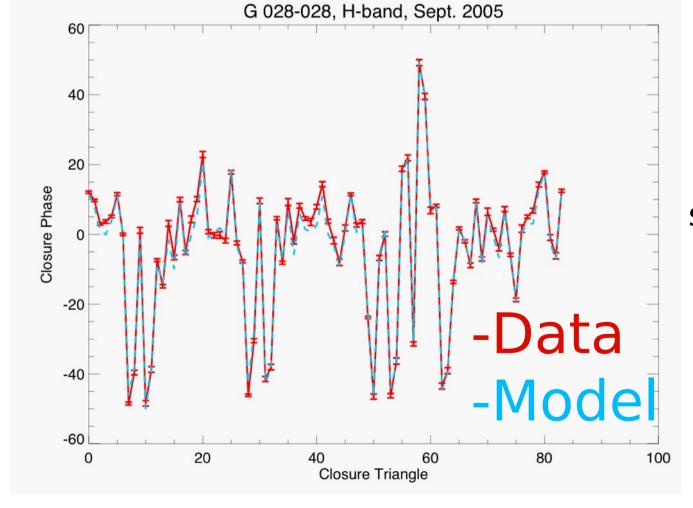
Closure phases are invariant to atmospheric phase

cancels out in closure phase sum

measured = intrinsic + atmospheric

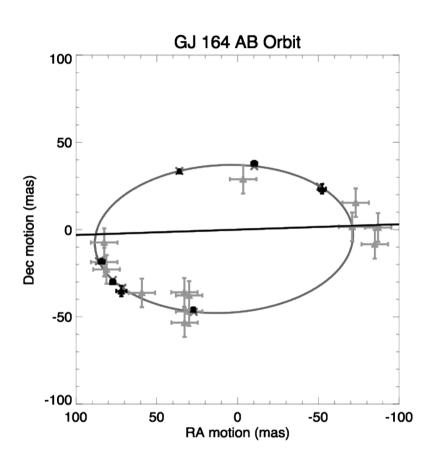
## **Binary systems**

3 parameters: angular separation, position angle, contrast Error estimate: closure phase scattering Small systematic error

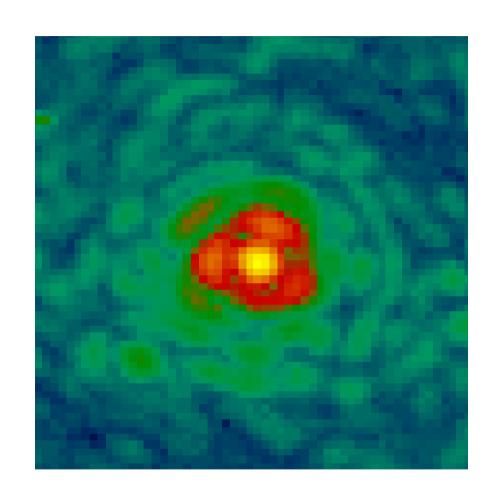


40 % strehl 0.3 deg scatter stability  $\sim \lambda/1000$  all passive!

#### An example of super-resolution



Black points: masking measurements Grey points: STEPS astrometry



H-band image of GJ164 by the Hale Telescope

 $\lambda/D = 66 \text{ mas}$ 

#### Beyond conventional aperture masking

Aperture masking offers high degree of calibration (= high precision) but is not very efficient (only a small fraction of the light is preserved). (u,v) coverage is limited by non-redundancy:

- large holes = high throughput, but fewer holes to avoid redundancy
- small holes = many holes, but small throughput
- → Aperture masking is most suitable for simple (compact) and bright objects

In aperture masking, non-redundancy needs to be imposed at the detector, not at the entrance aperture

→ it is possible to remap a dense grid of subapertures (redundant) into a sparse non-redundant array prior to Fizeau combination

# Example: the Fiber Imager foR Single Telescope (FIRST) instrument concept

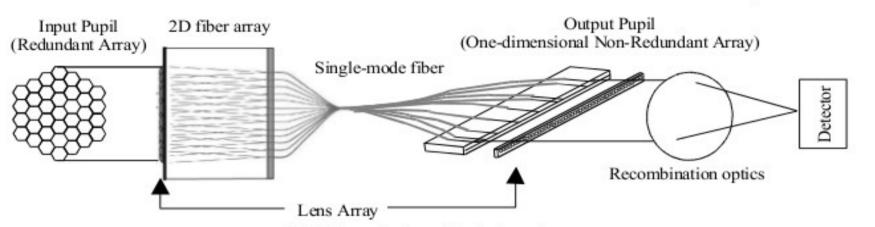
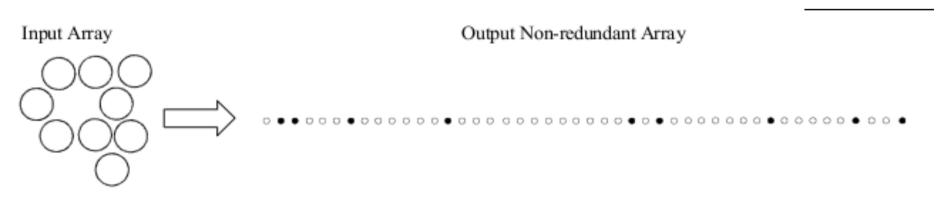


Fig.1. Schematic view of the insrtuemnts.

Remapping of a dense redundant array into a sparse non-redundant array Single mode fibers used for spatial filtering



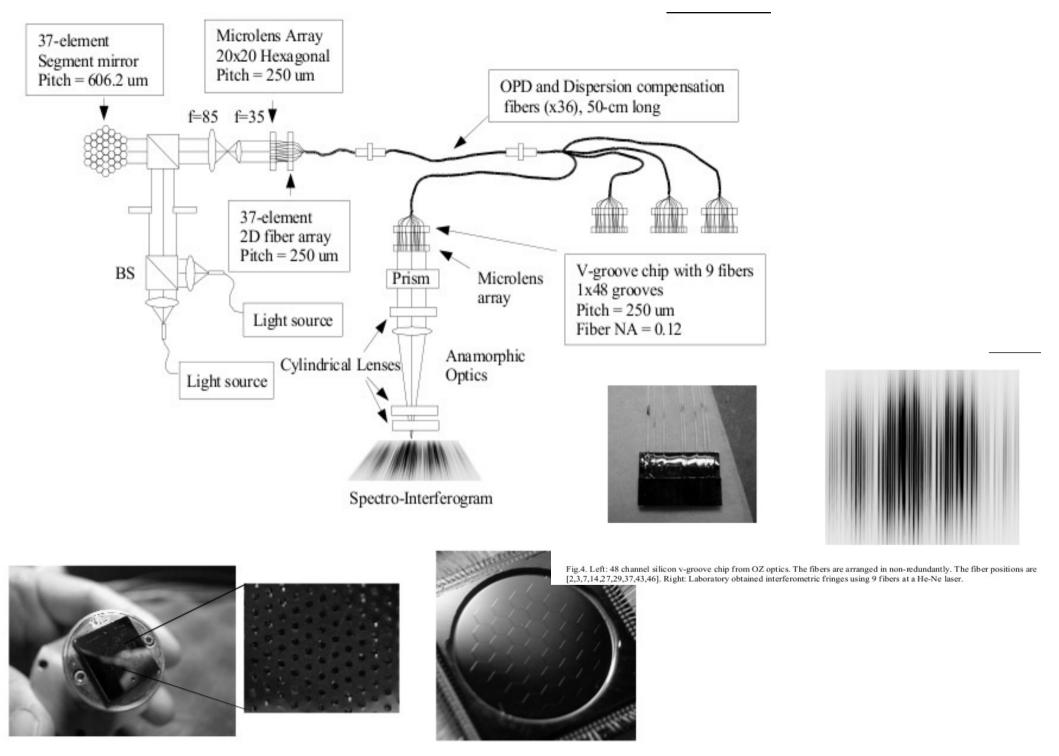
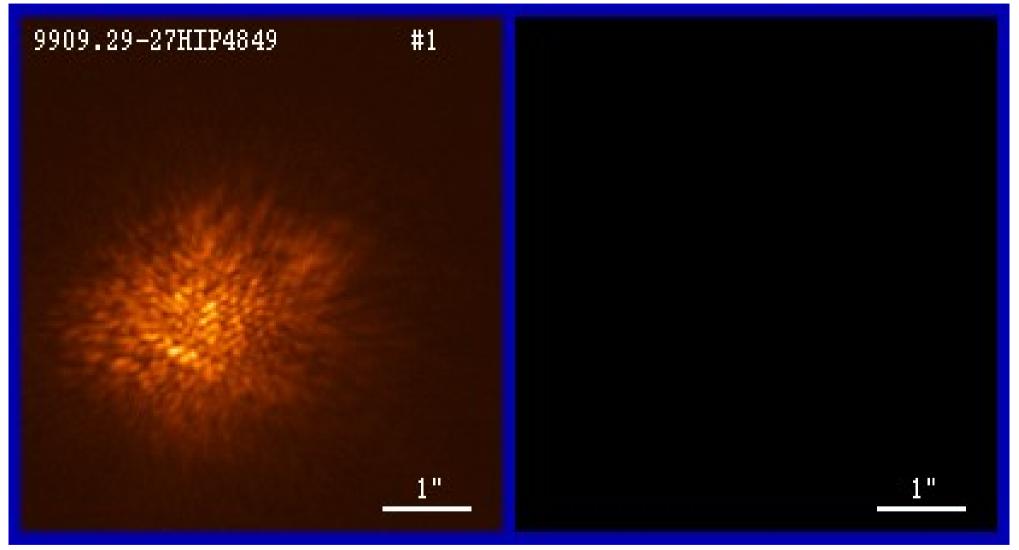


Fig.3. Left: 2D fiber array from FiberGuide Industries. The fiber pitch is 250  $\mu m$ . Right: Segmented Deformable Mirror from IRIS AO. The pitch between adjacent segments is 606.2  $\mu m$  including a 4  $\mu m$  gap.

# Speckle imaging



Real-time bispectrum speckle interferometry: 76 mas resolution. Frame rate of data recording and processing: ~ 2 frames per second. SAO 6 m telescope, K-band.

G. Weigelt, MPI for Radioastronomy, 1999

#### Speckle interferometry: reconstruction techniques

Single short exposure still contains signal at high spatial frequency.

Problem: phase and amplitude of the high spatial frequencies vary rapidly with time → long exposure will average high spatial frequencies to zero

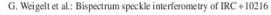
# Speckle interferometry (Power spectrum)

Solution developped by A. Labeyrie: average square modulus of images' Fourier transforms to obtain the square modulus FT of object

Problem: only amplitude of FT is recovered, not phase

# Speckle interferometry (Bispectrum)

Solution developped by K Weigelt: average bispectrum (equivalent of phase closures in sparse interferometry) to also recover phase



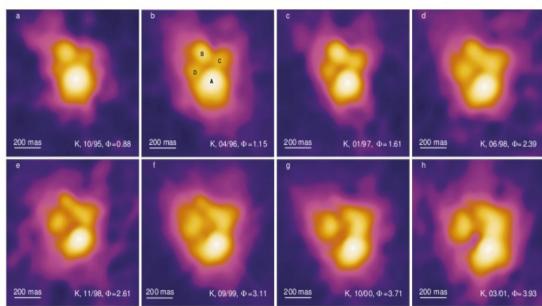


Fig. 1. K-band speckle reconstructions of IRC +10216 for 8 epochs from 1995 to 2001. The total area is 1"×1". All images are normalized to the brightest pixel and are presented with the same color table. North is up and east is to the left.

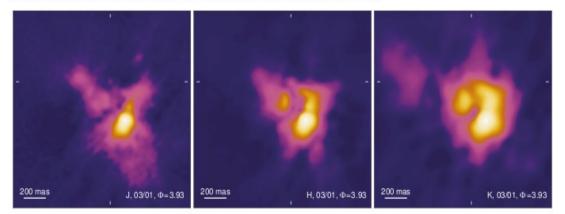


Fig. 2.  $J_{-}$ ,  $H_{-}$  and  $K_{-}$ band speckle reconstructions of IRC +10216 in March 2001. The total area is  $1.6" \times 1.6"$ . All images are normalized to the brightest pixel and are presented with the same color table. North is up and east is to the left. The tick marks indicate the likely position of the central star.

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#### **Speckle interferometry: reconstruction techniques**

Object image : f(x)

Its Fourier transform:  $F(u) = |F(u)| \exp[i\Phi(u)]$ 

Series of images (index n) is aquired. What is observed is the convolution of image and PSF:

$$s_n(x) = f(x) \circ psf_n(x)$$

In Fourier plane:

$$S_n(u) = F(u) \times PSF_n(u)$$

**Power spectrum speckle interferometry** (Labeyrie): measuring |F(u)|

$$<|S(u)|^2> = |F(u)|^2 < |PSF(u)|^2>$$

Speckle transfer function, calibrated on reference star

#### Bispectrum speckle interferometry (Weigelt):

Bispectrum : 
$$F^3(u_1,u_2) = F(u_1) F(u_2) F(-u_1-u_2) = |F^3(u_1,u_2)| \exp[i\psi(u_1,u_2)]$$

phase of object bispectrum

Average measured bispectrum:

$$<|S^3(u_1,u_2)|> = F^3(u_1,u_2) < |H^3(u_1,u_2)|>$$
 bispectrum transfer function, has zero phase

Allows recovery of bispectrum phase:  $\psi(u_1, u_2) = \Phi(u_1) + \Phi(u_2) - \Phi(u_1 + u_2)$ 

(This is a closure phase measurement)