

R. Kaiser¹, W. Guerin¹, G. Labeyrie¹, J.-P. Rivet², O. Lai², G. Labeyrie¹, S. Tolila¹, M. Hugbart¹, J.-G. Cuby³, Kevin Ho³, J. O'Meara⁴, A. Bouchez⁴, P. Wizinowich⁴, A. Adamson⁵, J. White⁵, E. Tapia⁵, N. Tamura⁶, O. Guyon⁶, J. Lozi⁶, M. Chun⁷, C. Dodds⁸, B. Allen⁷, D. Simons⁸

¹Institut de Physique de Nice, UCA, CNRS, France, ²Laboratoire Lagrange, UCA, OCA, CNRS, France, ³Canada-France-Hawaii Telescope, Waimea, Hawaii, USA, ⁴W.M. Keck Observatory, Waimea, Hawaii, USA, ⁵Gemini Observatory, Hilo, Hawaii, USA, ⁶Subaru Telescope, NAOJ, Hilo, Hawaii, USA, ⁷Institute for Astronomy, University of Hawaii, Hilo, Hawaii, USA, ⁸Institute for Astronomy, University of Hawaii, Manoa, Hawaii, USA

Introduction (1): Intensity correlations, the basics

Intensity correlation function:
$$g^{(2)}(\Delta \mathbf{r}, \tau) = \frac{\langle I(\mathbf{r}, t) I(\mathbf{r} + \Delta \mathbf{r}, t + \tau) \rangle}{\langle I(\mathbf{r}, t) \rangle^2}$$

In the spatial domain: $g^{(2)}(\mathbf{r}, \tau = 0)$

Classical picture: speckle from an extended disordered source.

The size of the speckle grain (= spatial correlation l_c) is related to the size of the source.

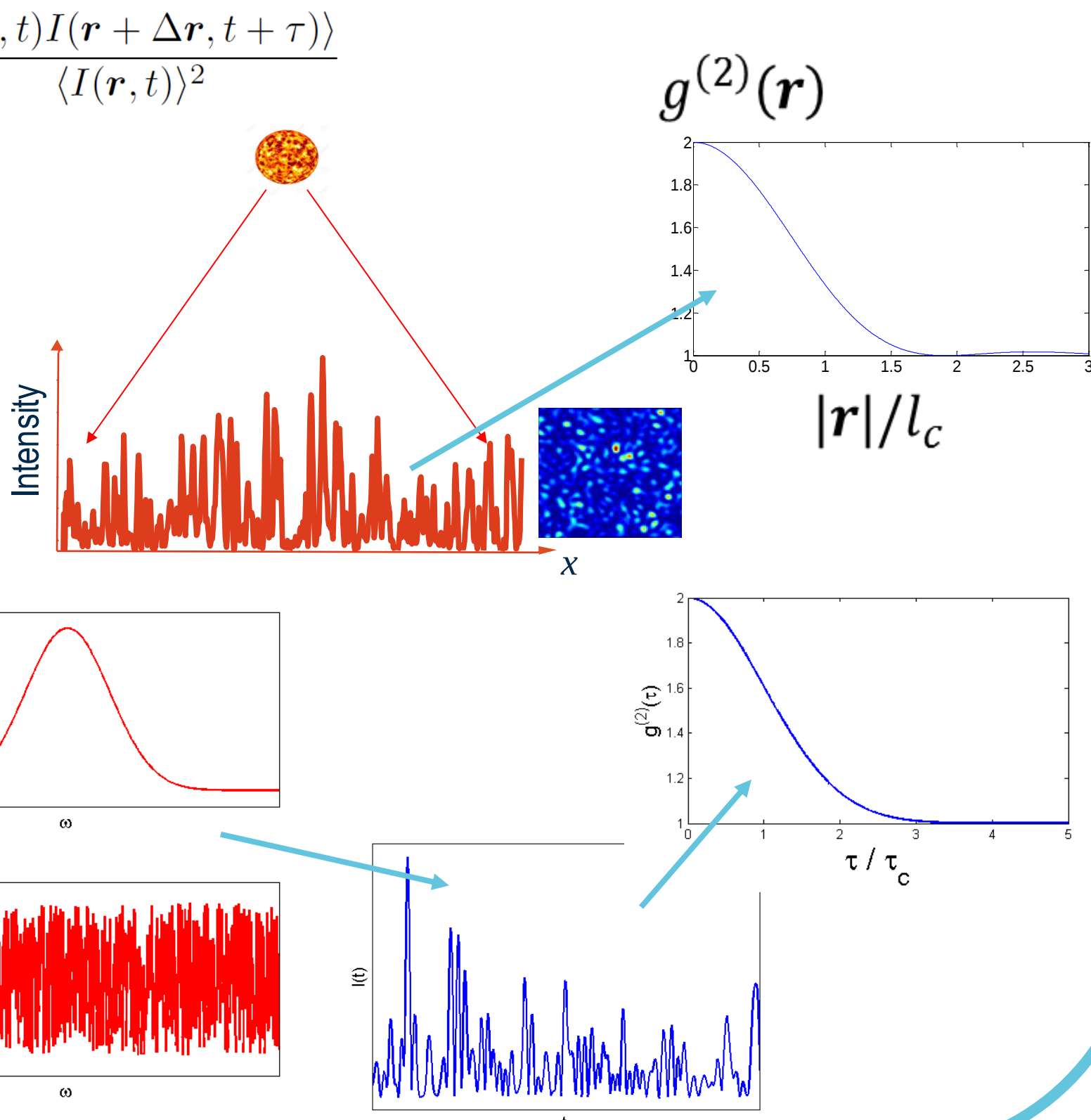
$$l_c = \lambda \cdot L/D$$

In the time domain: $g^{(2)}(\mathbf{r} = 0, \tau)$

Similar explanation: an incoherent (thermal) spectrum creates intensity noise (beating of the different frequency components).

The correlation time τ_c is related to the width of the spectrum.

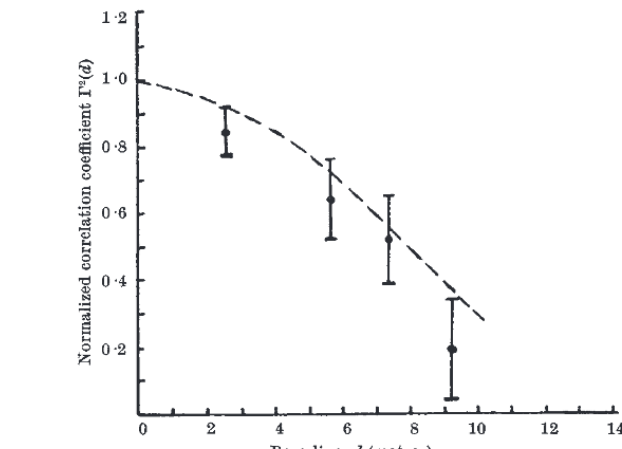
$$\tau_c = 1/\Delta\omega$$



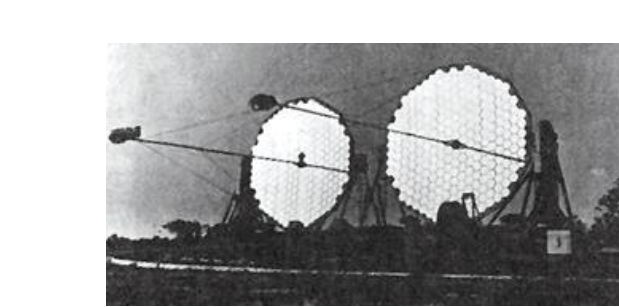
Introduction (2): Intensity correlations in astronomy

The past: the historical Hanbury Brown & Twiss experiments

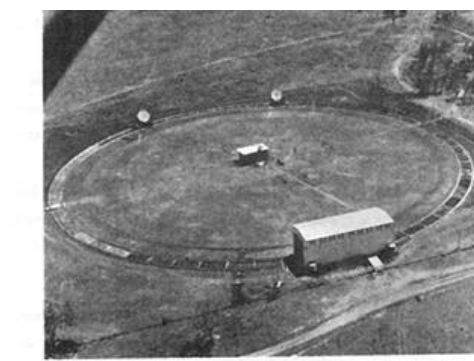
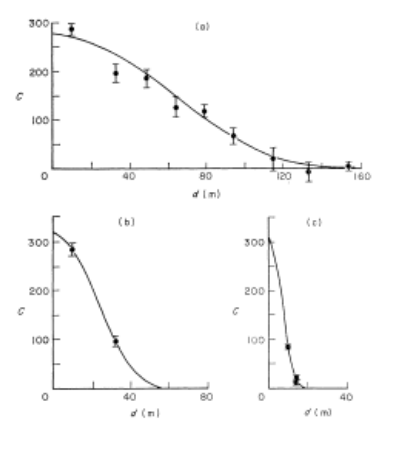
HBT 1956: First demonstration of $g^{(2)}(r)$ with 2 telescopes (measurement on Sirius) [1]



HB 1963-1974: The Narrabri stellar interferometer (Australia) [2,3]



Two collectors with $\varnothing = 6.5$ m. Measurement of the diameter of 32 stars with baselines between 10 and 188 m.



After 1975: the development of "direct interferometry" [4] put a halt to intensity interferometry. At the cost of combining the optical paths (with very precise and active delay lines), direct interferometry is much more efficient for the signal-to-noise ratio.

The future: towards a revival of intensity interferometry (see, e.g., [5,6])

Limitations of direct interferometry:

- Complicated and expensive
- Fringe tracking limits the accessible magnitude
- How to increase the baselines to km-scale ?

Intensity interferometry:

- Much simpler: no recombination, no delay line, insensitive to turbulence
- Separations can be increased to km or more
- Large array of many collectors possible (e.g. CTA) → high resolution imaging
- Need large collectors and very long integration times ☹️ but detectors and electronics made a lot of progress since the 1970s !!! → can be useful again

What are the performances and limitations of intensity interferometry with today's photonic technologies ?

$$SNR = \alpha N_{ph}(\lambda) A \sqrt{\frac{T_{obs}}{2\tau_{el}}}$$

α : detection efficiency → significant gain

N_{ph} : photon spectral flux

A : collection area

τ_{el} : electronic timing resolution → huge gain !

T : integration time

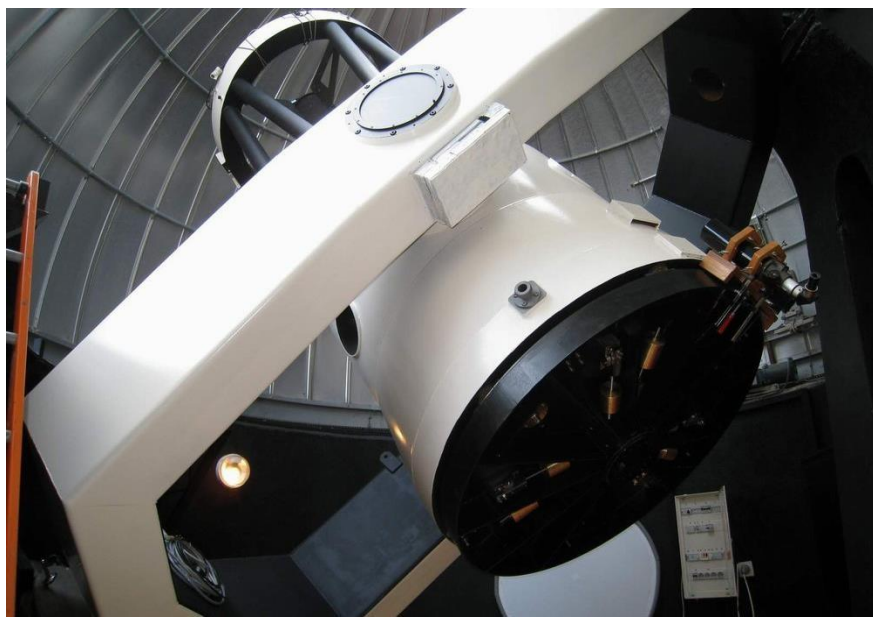
+ fiber technologies, digital electronics, time distribution at large distances...

And spectral multiplexing, $SNR \sim \sqrt{N}$, where N is the number of spectral channels;

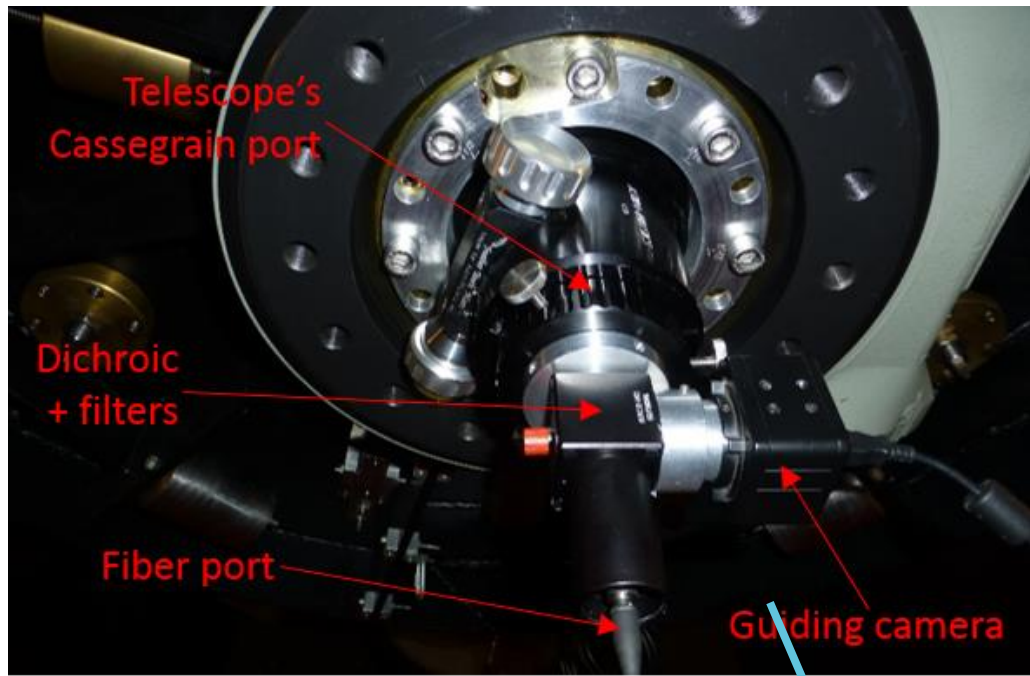
C2PU at Calern, OCA

Cassegrain telescopes:

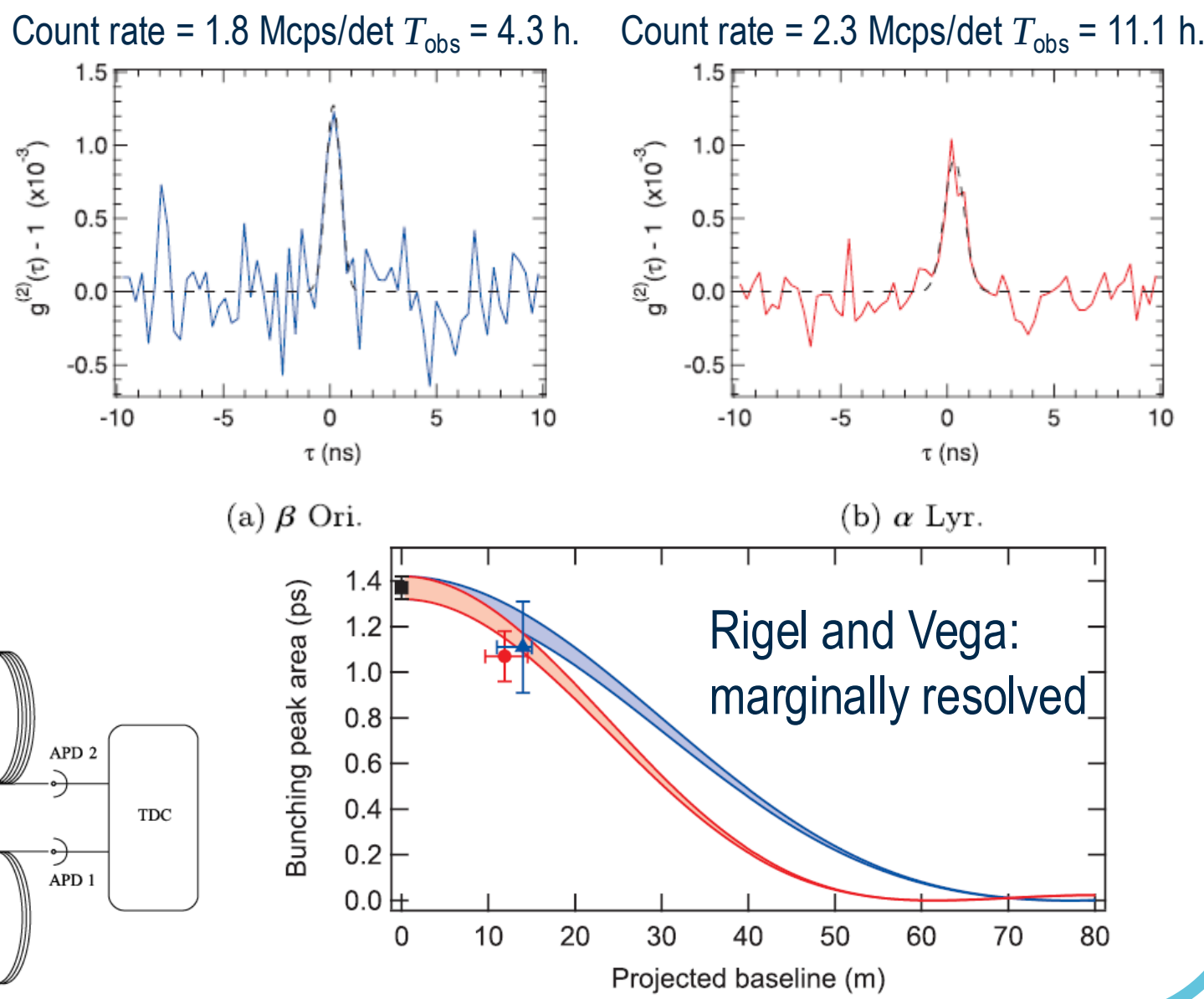
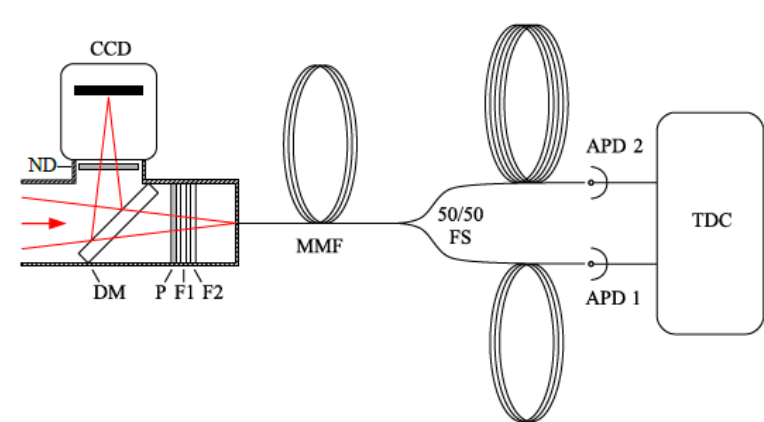
- $\varnothing = 1$ m
- Focal reducer → $f = 5.6$ m
- NA = 0.09 ; $f/5.6$
- PSF = 42 μ m for seeing = 1.5"



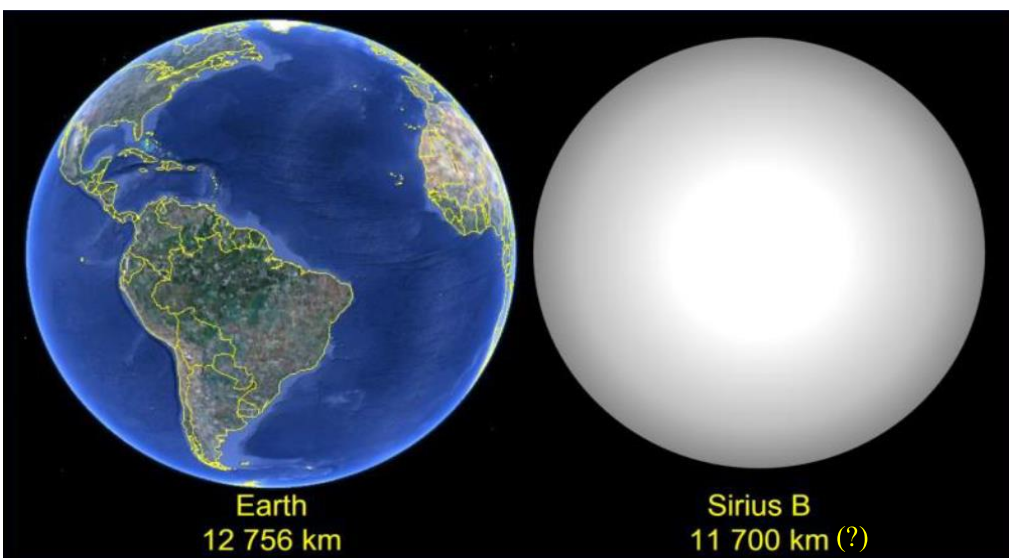
Results with starlight



We have developed a collaboration between quantum optics physicist and astronomers at Université Côte d'Azur since 2016 and have demonstrated the promise of the technique using our 1m telescopes to achieve original and meaningful astrophysical measurements [7, 8].



Angular resolution of a white dwarf



Count rates :
Sirius B
Quantum efficiency : 90%
Throughput : 20%

Keck: 110 000 cps
CFHT: 18 000 cps

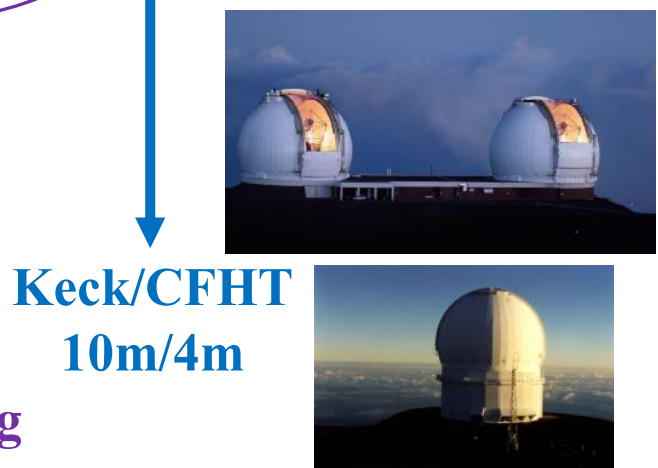
D=11700km
L=8.6 light years= 8 10¹⁶m
 $\Delta\theta=30\mu''$

These exotic objects are supported by Fermi electron degeneracy, and their diameter is estimated from their luminosity; such a measurement would uniquely constrain the diameter and demonstrate quantum mechanics at work on an astrophysical size object!

$$SNR = \sqrt{N_{channel}} A \eta F(\nu) |V(r)|^2 \sqrt{\frac{T_{obs}}{2\pi\tau_{el}}}$$



Multiplexing
N=16



Fast SPD
 $\tau_{el}=20ps$

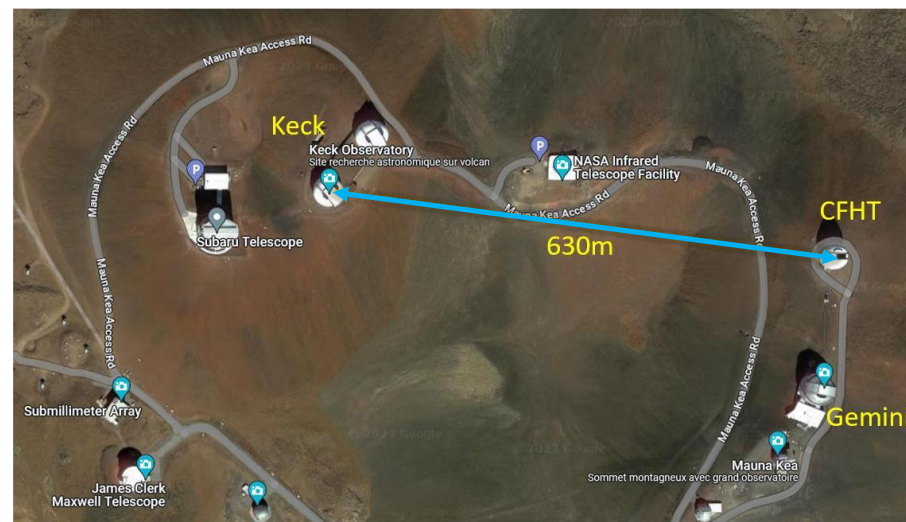
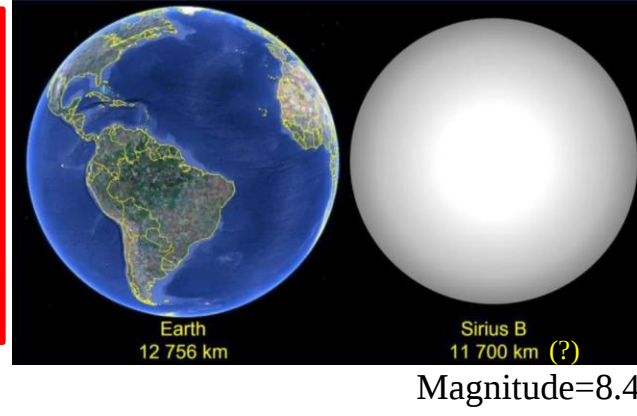
SNR : $\times 4 \times 40$

$\times 4 \Rightarrow \times 640$

$$T_{obs} \div 400\,000$$

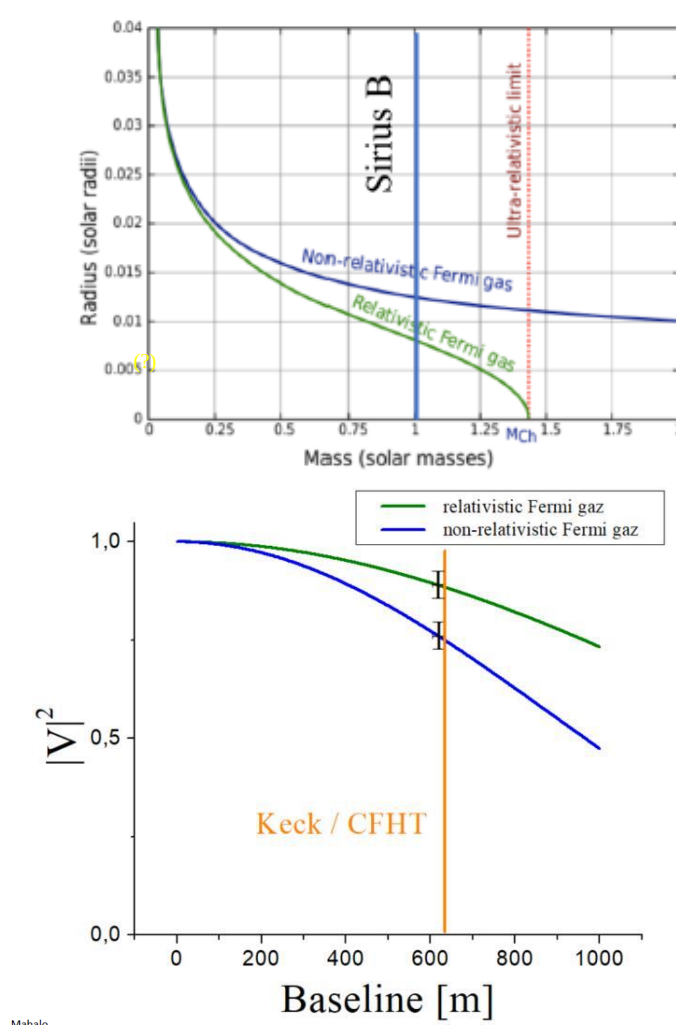
Path-opening on **Sirius B** (white dwarf) : quantum degenerate Fermi gas of electrons

SNR ≈ 6
in 1 hour
observation time !!!!
Beyond reach of present instruments



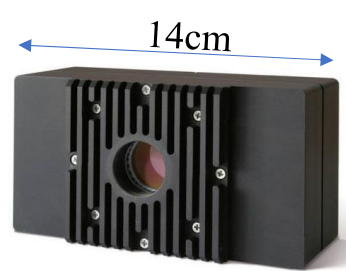
Mauna Kea @ Hawaii

Photon Bunching
• @ $\lambda = 420nm$
• D=630m

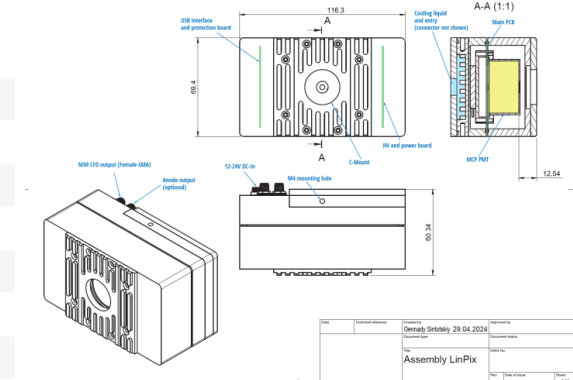


Technological progress driven by quantum optics and telecommunications has led to the development of new components (SPADs, SNSPDs) which have allowed to extend the sensitivity of the technique by orders of magnitude. We are developing various concepts for spectral multiplexing, such as an array of dichroic notch filters (right) to be used with LINPix detectors (below). We are also studying the possibility to use a linear SPAD array with a low time dispersion spectrograph

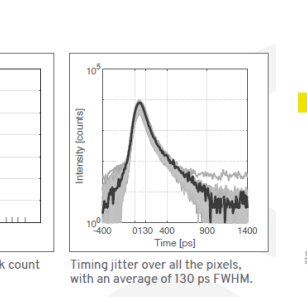
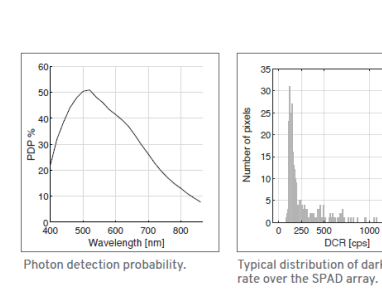
Photonscore : 2 x 16 LINPix



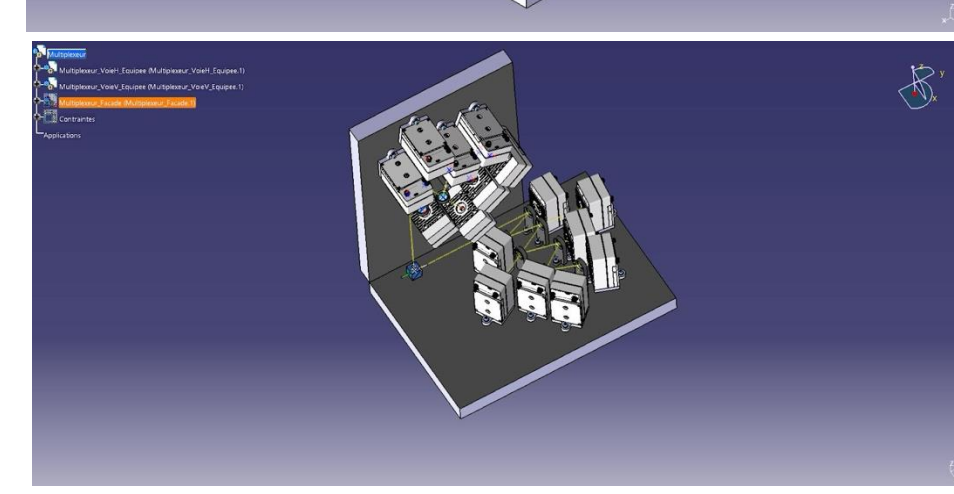
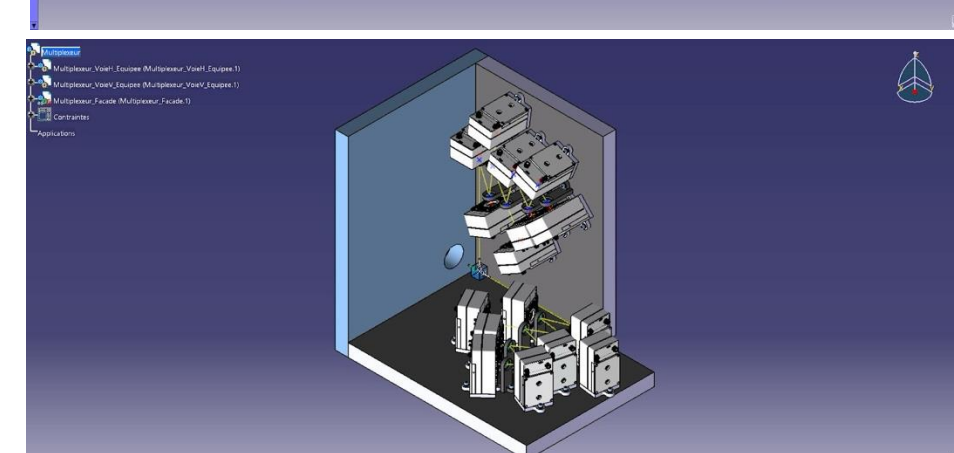
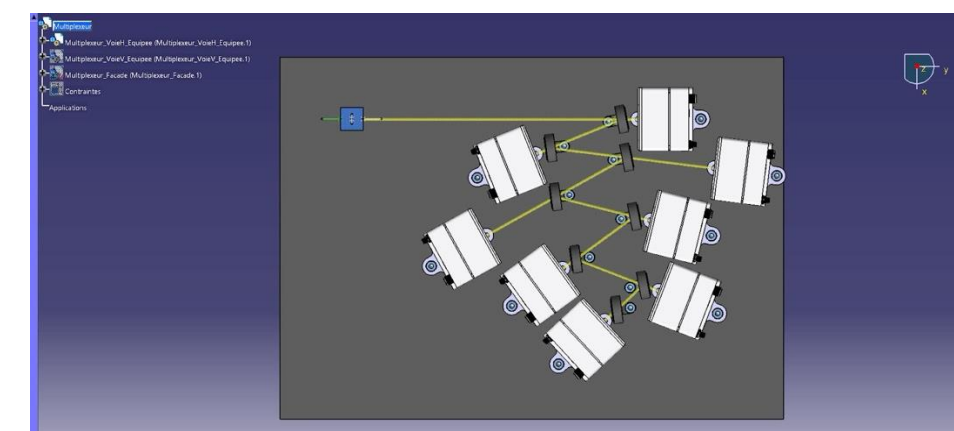
Parameter	Value
Max. recommended count rate (MHz)	100
Maximum count rate (MHz)	100
Dead time (ns)	10
Dark count rate (cps)	< 100
Timing jitter (ps FWHM)	< 100
Active area (mm)	40
Quantum yield (%)	> 20



Pi Imaging : 2 SPAD λ

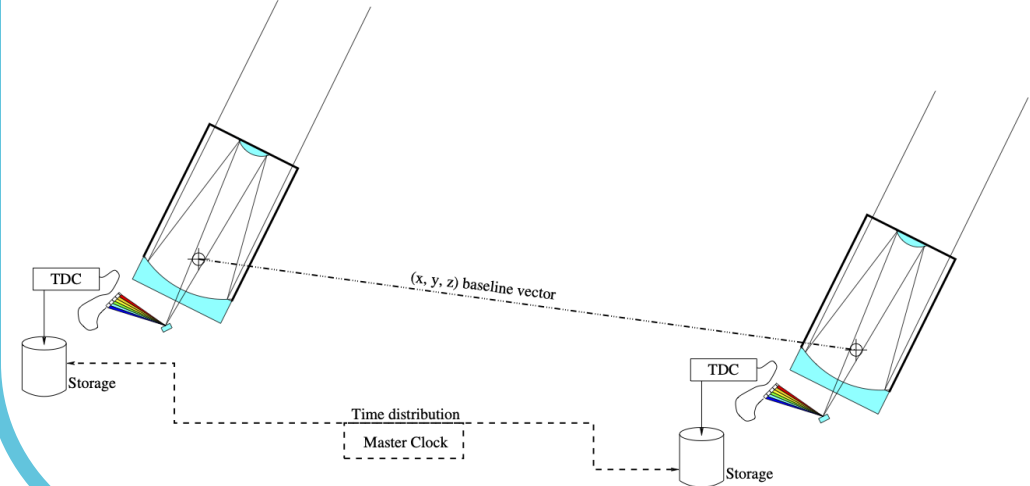


Parameter	Value
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Dark count rate (cps)	< 100
Timing jitter (ps FWHM)	< 100
Active area (mm)	40
Quantum yield (%)	> 20

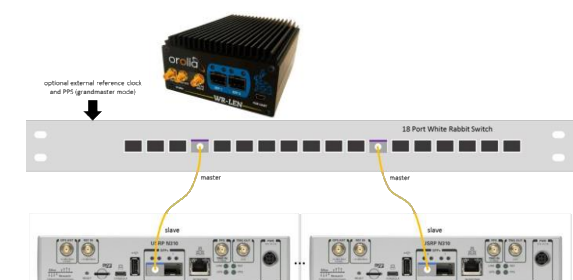


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Dark count rate (cps)	< 100
Timing jitter (ps FWHM)	< 100
Active area (mm)	40
Quantum yield (%)	> 20

1 ps



Photon arrival times will be recorded using a Time-to-Digital Converter (TDC) and written to disk for a posteriori correlations. However, clock at telescopes must be synchronised to picosecond level, and various solutions are available (below).



16 ps

References

[1] R. Hanbury Brown, R. Q. Twiss, Nature **178**, 1046 (1956).

[2] R. Hanbury Brown, Nature **218**, 637 (1968).

[3] R. Hanbury Brown, J. Davis, L. R. Allen, MNRAS **167**, 121 (1974).

[4] A. Labeyrie, ApJ **196**, L71 (1975).

[5] D. Dravins, Proc. SPIE **9907**, 0M (2016).

[6] J.-P. Rivet *et al.*, Exp. Astron. (2018).

[7] W. Guerin *et al.*, MNRAS **472**, 4126 (2017).

[8] W. Guerin *et al.*, MNRAS **480**, 245 (2018).

[9] J.M. Mariotti *et al.*, A&AS, **116**, 381 (1996)

[10] G. Perrin *et al.*, Science, **311**, 194 (2006)

[11] O. Lai *et al.*, SPIE Proc. **4838**, 1296 (2003)